Localization of Inner-Shell Photoelectron Emission and Interatomic Coulombic Decay in Ne$_2$

K. Kreidi[1,2], T. Jahnke[1], Th. Weber[3], T. Havermeier[1], R. E. Grisenti[1,4], X. Liu[5], Y. Morisita[6], S. Schössler[1], L. Ph. H. Schmidt[1], M. Schöffler[1], M. Odenweller[1], N. Neumann[1], L. Foucar[1], J. Titze[1], B. Ulrich[1], F. Sturm[1], C. Stuck[1], R. Wallauer[1], S. Voss[1], I. Lauter[1], H. K. Kim[1], M. Rudloff[1], H. Fukuzawa[5], G. Prümper[5], N. Saito[6], K. Ueda[5], A. Czasch[1], O. Jagutzki[1], H. Schmidt-Böcking[1], S. K. Semenov[7], N. A. Cherepkov[7] and R. Dörner[1*]

[1] Institut für Kernphysik, J. W. Goethe Universität, Max-von-Laue-Str. 1, 60438 Frankfurt, Germany
[2] DESY, Notkestrasse 85, 22607 Hamburg, Germany
[3] Lawrence Berkeley National Laboratory, Berkeley CA 94720, USA
[4] Gesellschaft für Schwerionenforschung, Planckstr. 1, 64291 Darmstadt, Germany
[5] Institute of Multidisciplinary Research for Advanced Materials, Tohoku University, Sendai 980-8577, Japan
[6] National Metrology Institute of Japan, AIST, Tsukuba 305-8568, Japan
[7] State University of Aerospace Instrumentation, 190000 St. Petersburg, Russia

We used Cold Target Recoil Ion Momentum Spectroscopy (COLTRIMS) to investigate the decay of Ne$_2$ after K-shell photoionization. The breakup into Ne^{1+} / Ne^{2+} shows interatomic Coulombic decay (ICD) occurring after a preceding atomic Auger decay. The molecular frame angular distributions of the photoelectron and the ICD electron show distinct, asymmetric features, which imply localization of the K-vacancy created at one of the two atomic sites of the Ne$_2$ and an emission of the ICD electron from a localized site. The experimental results are supported by calculations in frozen core Hartree-Fock approach.

PACS numbers:

Are inner-shell holes in homonuclear diatomic molecules localized at one of the atoms or delocalized over the two equivalent sites? This question has been discussed highly controversial in literature for more than 35 years now (see e.g. [1, 2, 3, 4, 5, 6, 7, 8, 9, 10, 11] and our discussion below). Here we report on an experiment answering that question for the Ne$_2$ Van-der-Waals-molecule. Intuitively strong arguments for both opinions may be found: the K-shell wave function is very tightly confined to the nuclei and the overlap between inner-shell orbitals at different atoms of a molecule is usually negligibly small [12]. Hence, the geometry of the problem suggests to think of individual localized atomic wave functions. The symmetry of the problem however suggests the opposite: both sites of the diatomic molecule are indistinguishable and therefore the total molecular wave function has to have well defined *gerade* (g) or *ungerade* (u) symmetry. In order to construct the molecular many-body wave function it seems natural to employ only symmetry adapted single electron orbitals. A core level hole would then have well defined g or u symmetry and hence be delocalized over the two sites. This approach is used in today's state of the art theory concerning the photoionization and decay of Ne$_2$ [13].

To address this question of a possible localization in a quantum mechanically meaningful and experimentally accessible way one has to relate it to a measurable observable. Prime candidate is the energy of the state measured by photoelectron spectroscopy. Theoretical claims for core hole localization based on the energy date back to classic works of Snyder [1] and Bagus [2]. They showed that allowing symmetry broken - i.e. localized - basis states for inner-shell vacancies in a Hartree-Fock calculation lowers the energy and yields better agreement with the experiment. It was later argued that this is a peculiarity of the Hartree-Fock approach [3]. More sophisticated calculations on N$_2$ today very well reproduce the experimentally observed energy splitting of about 100 meV between the $1\sigma_g$ and $1\sigma_u$ core-ionized states [4, 5]. This is generally taken as evidence for the delocalized character of the inner-shell hole.

As an alternative observable being sensitive to core hole localization the angular distribution of the photoelectrons in the laboratory system was suggested [6]. Corresponding experiments on nitrogen molecules show good agreement with calculations using delocalized orbitals [7]. This conclusion is also supported by resonant soft X-ray emission experiments on O$_2$ [8]. There the parity selection rules indicate the symmetry and with it the delocalization of the core-excited state. In contrast to that partial localization is shown in experiments on singly substituted ^{14,15}N$_2$ [9]. Furthermore, experiments exciting an inner-shell electron inducing fast dissociation and observing the Auger electron emitted during this dissociation show a clear Doppler-shift in the Auger peak. This proves that they are emitted from a localized source flying towards to or away from the observer [10, 14].

In the present work we use an even more sophisticated probe for localization which is the electron angular distribution in the *body-fixed frame* of the molecule [14, 15, 16]:
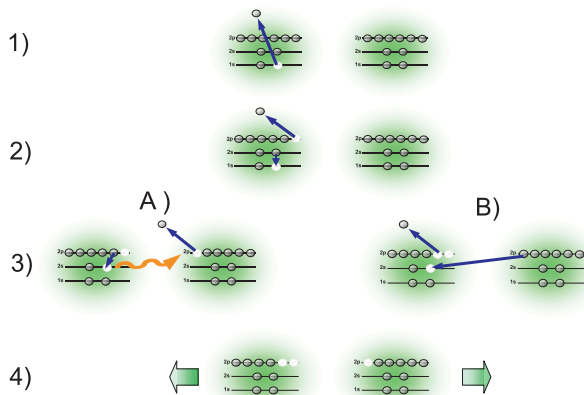
FIG. 1: $1s$ photoionization (1) leads to an Auger decay resulting in a dicationic one-site state of Ne_2 (2). The Auger decay is followed by interatomic Coulombic decay (3) where the ion's excitation energy is emitted via a virtual photon transfer (A) or via electron transfer (B). Ending up in a state where both atoms are charged positively, the dimer fragments back-to-back in a Coulomb explosion (4).

we investigate the angular distribution of a photoelectron and an electron emitted via interatomic Coulombic decay (ICD) [17] following Auger decay [18] for fixed molecular orientations. If the electron is emitted from a delocalized source its angular distribution will be symmetric with respect to the two atoms of the molecule (see e.g. [19, 20]). If, however, it is emitted from one site, then the angular distribution can show a strong asymmetry due to interference of the electron-wave that is multiply scattered in the molecular potential. For example, a strong forward focussing by the atomic neighbor can occur [21]. In a homonuclear diatomic molecule such an asymmetry however can only be observed for an asymmetric breakup of the molecule into two fragments of different charge which make the two ends of the molecule experimentally distinguishable. A first experiment along this line on N_2 found no asymmetry with respect to the N^{1+}/N^{2+} fragments [20]. Here we investigate K-shell ionization of Ne_2 instead of N_2 for several reasons. Firstly, K-shell ionized Ne_2 fragments into the asymmetric breakup channel Ne^{1+} / Ne^{2+}. Secondly the K-shell radius of Ne_2 is 30 times smaller than the internuclear distance R of 5.86 a.u. The bond is purely Van-der-Waals, the binding energy is only 3 meV and the g/u splitting is negligibly small compared to the natural line width. All this might make it plausible to think of Ne_2 as two Neon atoms sitting close together with their K-shells being independent of each other. Thirdly however the opposing symmetry argument in favor of a delocalized description of all electrons by symmetry adapted wave functions of well defined g or u parity is as valid for Ne_2 as it is for N_2. Hence the key arguments keeping the question of possible localiza-

tion open today in covalently bound molecules hold for Van-der-Waals dimers, as well.

Santra *et al.* [18] theoretically investigated the decay of Ne_2 after $1s$ ionization. They suggested that one-site Auger final states $Ne^{2+}(2s^{-1}2p^{-1})$ / Ne decay further to a two-site state $Ne^{2+}(2p^{-2})$ / $Ne^{1+}(2p^{-1})$ via ICD. Fig. 1 shows a sketch of the different pathways of decay leading to this set of reaction products. Both decay routes start with K-shell ionization. In a localized scenario it is followed by an Auger decay of the same atom of the dimer yielding a vacancy in an inner valence shell ($Ne^{2+}(2s^{-1}2p^{-1})$ / Ne). Now two different channels of ICD may occur depending on the parity of the doubly charged states [22, 23]. Pathway A) shows ICD via virtual photon exchange, B) depicts the competing process involving the transfer of an electron. As both A) and B) lead to a triply charged species, Ne_2 fragments in a Coulomb explosion as a final step. The existence of ICD with its two center nature is the key feature in Ne_2, making it possible to trace core hole localization through all steps of the decay of Ne_2. The sketch of the decay in Fig. 1 shows fully localized electrons and holes. Thus it may seem that the pure existence of ICD already proofs a localization. This is not true however. While most discussions on ICD on an intuitive level use figures similar to our Fig. 1, all actual calculations on ICD, including the ones performed by Santra for the present system [18], work with molecular wave functions of well defined g or u symmetry, i.e. they assume complete delocalization. The ICD calculations using symmetry adapted wave functions yield excellent agreement of the experimentally observed energies of the states and of the energy of the ICD electrons [13, 24]. These calculations report ICD electron energies separately for g and u symmetry again highlighting the delocalized character of the calculations [13].

The experiment was performed at beamline UE56/1-SGM of the Berlin Synchrotron (BESSY) using the COLTRIMS technique [25]. As a source for Ne_2 a supersonic jet that was precooled to 160 K was employed. It was crossed with the photon beam. Products from the photoreaction were guided by an electric field of 20 V/cm and a magnetic field of 6 Gauss towards two channel plate detectors with delayline readout [26]. Electrons up to 12 eV and ions up to 10 eV were detectable with a solid angle of 4π.

Fig. 2 shows our experimental results for the measured electron energies and the sum of the kinetic energies of the ionic fragments (KER, kinetic energy release). With a photon energy of $h\nu = 881.2$ eV, events located at an electron energy of 11 eV correspond to photoelectrons from Ne $1s$ ionization. Electrons originating from ICD are identified as diagonal lines A) and B). As the sum of the energy of the ionic fragments and the ICD electron is a constant, these diagonal lines with a slope of $-45°$ are a clear evidence for ICD as shown in [24].
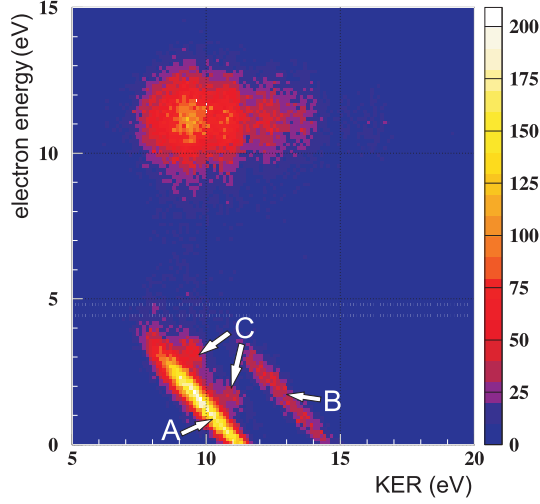
The ICD channels, labeled A) and B) in Fig. 2,

FIG. 2: The kinetic energy of the ionic fragments in dependence of the electron energy. The photon energy is $h\nu = 881.2$ eV resulting in an energy of the $1s$ photoelectron of 11 eV. The two diagonal lines are a clear evidence for interatomic Coulombic decay (ICD). Line A) is produced via the virtual photon exchange yielding a sum energy of 11.1 eV. Line B) (with a sum energy of 14.3 eV) occurs as ICD via electron transfer happens. The origin of channel C), showing two separate islands, is not quite clear.
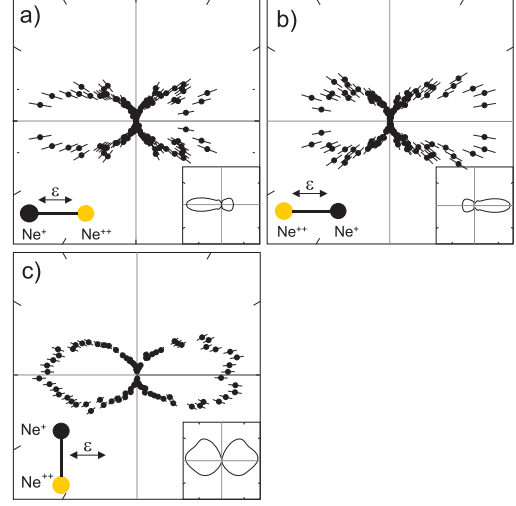

FIG. 3: Angular distribution of the 11 eV $1s$ photoelectron in dependence of the orientation of the dimer axis and the direction of the polarization vector ε (horizontal). In a) and b) the dimer is aligned parallelly to the polarization vector within an angle of $\pm$ 3°. In c) the plane is defined by the dimer axis and the polarization of the light. The dimer is aligned perpendicularly to the polarization vector and the photoelectron is fixed relative to the plane within an angle of $\pm$ 10°. The asymmetry, being a result of core hole localization, is clearly visible in the experimental data (circles). Solid line: frozen core Hartree-Fock calculation assuming a localized emission.

have a different KER which shows that the decay occurs at different internuclear distances. Similar to the findings in [23] both channels are created in an IC decay with the symmetry of the involved states and the difference in kinetic energy of the ions implying the occurrence of the two different contributions to ICD: channel A) with a sum energy of 11.1 eV represents the decay from the inner valence excited one-site state Ne^{2+} $(2s^{-1}2p^{-1})$ $[^1P]$ / Ne $[^1S]$ to the two-site state Ne^{2+} $(2p^{-2})$ $[^1D]$ / Ne^{1+} $(2p^{-1})$ $[^2P]$. This so called 'direct' IC decay happens via an exchange of a virtual photon as sketched as pathway A in Fig. 1. The measured KER of $\sim$ 8 eV to $\sim$ 11 eV corresponds to an internuclear distance of $\sim$ 3.6 Å to $\sim$ 2.6 Å [27]. The mean value of this range is close to 3.05 Å which is the distance of the neon atoms in the ground state of Ne_2 and thus consistent with corresponding results in [23]. Channel B) with a sum energy of KER + electron energy = 14.3 eV occurs at much higher KER ($\sim$ 11 eV to $\sim$ 14 eV) which is equivalent to a internuclear distance of only $\sim$ 2.6 Å to $\sim$ 2.0 Å. Here the one-site state Ne^{2+} $(2s^{-1}2p^{-1})$ $[^1P]$ / Ne $[^1S]$ decays into the two-site state Ne^{2+} $(2p^{-2})$ $[^3P]$ / $Ne^{1+}(2p^{-1})$ $[^2P]$ after electron transfer, as shown by pathway B in Fig. 1. The decay is described by the 'exchange' part of the electron-electron Coulomb matrix element. In this case the spatial overlap of the involved wave functions is the crucial contribution to the decay probability. Therefore pathway B is sup-

pressed at large internuclear distances where pathway A is still open [23]. Beside channels A and B our experimental results also show two small islands, labeled as C in Fig. 2. The sum energy of KER + electron energy = 12.1 eV indicates the decay from the dicationic state Ne^{2+} $(2s^{-2})$ $[^1S]$ / Ne $[^1S]$ to the two-site state Ne^{2+} $(2s^{-1})(2p^{-1})$ $[^3P]$ / Ne^{1+} $(2p^{-1})$ $[^2P]$. The KER of this decay is $\sim$ 11 eV which is equivalent to $R = 2.6$ Å. This channel also represents a spin-flip ICD where electron transfer may play a role but until now it is not clear why it shows two separate islands instead of a complete diagonal line.

Fig. 3 shows the angular distribution of the emitted photoelectron in the body fixed frame of Ne_2. Clearly, the photoelectron angular distribution is asymmetric - the photoelectron is preferably emitted towards the singly charged dimer fragment. To validate our findings Fig. 3(a) and 3(b) show the results where the doubly charged ion is emitted in two opposite directions within the laboratory frame. As a support for our interpretation we compared the experimental angular distributions with a theoretical prediction for the case of a completely localized electron. The calculation was performed within the Hartree-Fock approximation using the method described
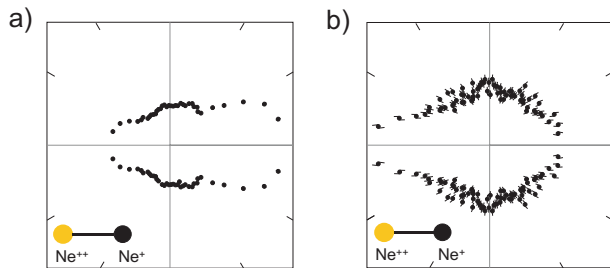
FIG. 4: Angular distribution of the ICD-electron in the dimer frame. The dimer is aligned horizontal and it is integrated over the orientation of the polarization vector. a) ICD electrons which are created by virtual photon exchange (Channel A) in Fig.1), b) angular distribution of the ICD electrons emitted after transfer of an electron (Channel B in Fig.1). As suggested by the sketch in Fig.1 the direction of the asymmetry switches from a) to b) depending on the decay path.

in [28]: at first the ground state of Ne_2 was computed in order to obtain the initial $1\sigma_g$ and $1\sigma_u$ wave functions. The photoelectron wave function was calculated in the frozen core Hartree-Fock approximation. Localization of the initial hole after photoionization on the right or left atom ($|r\rangle$ or $|l\rangle$ states) was introduced by taking a linear combination of the symmetry adapted $1\sigma_g$ and $1\sigma_u$ wave functions, $|r\rangle = (|1\sigma_g\rangle + |1\sigma_u\rangle)/\sqrt{2}$ and $|l\rangle = (|1\sigma_g\rangle - |1\sigma_u\rangle)/\sqrt{2}$. The full line down right in Fig. 3(a), 3(b) and 3(c) shows the prediction for the assumption of fully localized emission of the photoelectron from the doubly charged part of Ne_2. It nicely resembles the asymmetric shape of the measured distribution, overestimating the total magnitude of the asymmetry. Fig. 3(c) shows the same distribution with the dimer being oriented perpendicularly to the polarization axis of the photon beam. As expected no left/right asymmetry is visible in that case, but a small up/down asymmetry, which is again overestimated by our calculations that imply complete localization.

As a next step we investigate the corresponding angular distribution for the ICD-electron. Fig. 4(a) shows the distribution for decay path A, Fig. 4(b) the one for decay channel B. Both angular distributions are - just as the one for the photoelectron - asymmetric with respect to the two atomic centers of the dimer. A striking difference is visible for the preferred direction of emission. While the ICD electron is emitted towards the Ne^{1+} for channel A) it is emitted towards the Ne^{2+} in case of pathway B). This implies the ICD electron being emitted from opposite sites in the two cases, which is exactly what one would intuitively expect from comparing the two processes in Fig. 1. This finding is maybe even more surprising, as the emitted ICD-electron originates from a valence shell: for an excited/ionized Van-der-Waals molecule valence electrons are often viewed as

delocalized molecular orbitals.

In conclusion, the molecular frame angular distribution of photoelectrons and electrons from interatomic Coulombic decay are found to be strongly asymmetric for asymmetric breakup channels of Ne_2. For the photoelectron this finding directly proves that a photon induced core hole in Ne_2 is best thought of as being localized. The observed asymmetry for the ICD electrons shows that in addition to the core hole also the 2s hole created by the Auger decay and the valence orbital that emits the ICD electron show strong features of localization.

* Electronic address: doerner@atom.uni-frankfurt.de
[1] L. C. Snyder, J. Chem. Phys. **55**, 95 (1971).
[2] P. S. Bagus, J. Chem. Phys. **56**, 224 (1972).
[3] J. A. Kintop, W. V. M. Machado, and L. G. Ferreira, Phys. Rev. A **43**, 3348 (1991).
[4] A. Thiel, J. Schirmer, and H. Köppel, J. Chem. Phys. **119**, 2088 (2003).
[5] M. Ehara et al, J. Chem. Phys. **124**, 124311 (2006).
[6] D. Dill et al, Phys. Rev. Lett. **41**, 1230 (1978).
[7] M. Matsumoto et al, J. Phys. B **39**, 375 (2006).
[8] P. Glans et al, Phys. Rev. Lett. **76**, 2448 (1996).
[9] D. Rolles et al, Nature **437**, 711 (2005).
[10] O. Björneholm et al, Phys. Rev. Lett. **84**, 2826 (2000).
[11] C. T. Chen, Y. Ma, and F. Sette, Phys. Rev. A **40**, 6737 (1989).
[12] N. Kosugi, Chem. Phys. **289**, 117 (2003).
[13] S. Scheit et al, J. Chem. Phys. **121**, 8393 (2004).
[14] A.V. Golovin et al, Phys. Rev. Lett. **79**, 4554 (1997).
[15] N. Saito et al, J. Phys. B **38**, L277 (2005).
[16] J. Adachi et al, J. Phys. B **40**,F285 (2007).
[17] L. S. Cederbaum, J. Zobeley, and F. Tarantelli, Phys. Rev. Lett. **79**, 4778 (1997).
[18] R. Santra, and L. S. Cederbaum, Phys. Rev. Lett. **90**, 153401 (2003). R. Santra, and L. S. Cederbaum, Phys. Rev. Lett. **94**, 199901(E)(2005).
[19] A. A. Pavlychev et al, Phys. Rev. Lett. **81**, 3623 (1998).
[20] T. Weber et al, J. Phys. B **34**, 3669 (2001).
[21] H. C. Poon, and S. Y. Tong, Phys. Rev. B **30**, 6211 (1984).
[22] R. Santra, J. Zobeley, and L. S. Cederbaum, Phys. Rev. B **64**, 245104 (2001).
[23] T. Jahnke et al, Phys. Rev. Lett. **99**, 153401 (2007).
[24] T. Jahnke et al, Phys. Rev. Lett. **93**, 083002 (2004).
[25] J. Ullrich et al, Rep. Prog. Phys. **66**, 1463 (2003).
[26] O. Jagutzki et al, Nucl. Instr. Meth. A **477**, 244 (2002).
[27] T. Weber et al, Nature **431**, 437 (2004).
[28] S. K. Semenov et al, Phys. Rev. A **61**, 032704 (2000).